

1. What are the eigenvalues of S^2 and S_z for the spin function

$$(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3))/\sqrt{3} \quad \text{** normalization function wrong in question}$$

S	M_S	Spin Adapted Configuration
3/2	+ 3/2	$^4\Phi_{3/2} = \alpha\alpha\alpha$
3/2	+ 1/2	$^4\Phi_{1/2} = 1/3^{1/2} (\alpha\beta\alpha + \beta\alpha\alpha + \alpha\alpha\beta)$
3/2	- 1/2	$^4\Phi_{-1/2} = 1/3^{1/2} (\alpha\beta\beta + \beta\alpha\beta + \beta\beta\alpha)$
3/2	- 3/2	$^4\Phi_{-3/2} = \beta\beta\beta$
1/2	+ 1/2	$^2\Phi_{1/2} = 1/6^{1/2} (\alpha\beta\alpha + \beta\alpha\alpha - 2\alpha\alpha\beta)$
1/2	+ 1/2	$^2\Phi_{1/2} = 1/2^{1/2} (\beta\alpha\alpha - \alpha\beta\alpha)$
1/2	- 1/2	$^2\Phi_{-1/2} = 1/6^{1/2} (\beta\alpha\beta + \beta\beta\alpha - 2\alpha\beta\beta)$
1/2	- 1/2	$^2\Phi_{-1/2} = 1/2^{1/2} (\beta\beta\alpha - \beta\alpha\beta)$

are the **set of spin functions for 3 electrons which are in separate space orbitals** (e.g. $1s^1 2s^1 2p^1$ configuration of excited Li)

The goal of the problem is to show that the values of S , M_S for the $^4\Phi_{1/2}$ spin state are $(3/2, +1/2)$
 i.e. $S^2 \Phi_{1/2} = S(S+1) \Phi_{1/2} = (3/2)(5/2) \Phi_{1/2} = 15/4 \Phi_{1/2}$ and
 $S_z \Phi_{1/2} = M_S \Phi_{1/2} = (+1/2) \Phi_{1/2}$

$$\begin{aligned} S^2 &= (S_1 + S_2 + S_3) \cdot (S_1 + S_2 + S_3) = S_1^2 + S_2^2 + S_3^2 + 2(S_1 \cdot S_2 + S_1 \cdot S_3 + S_2 \cdot S_3) \\ &= S_1^2 + S_2^2 + S_3^2 + 2(S_{1x} \cdot S_{2x} + S_{1y} \cdot S_{2y} + S_{1z} \cdot S_{2z} + S_{1x} \cdot S_{3x} + S_{1y} \cdot S_{3y} + S_{1z} \cdot S_{3z} \\ &\quad + S_{2x} \cdot S_{3x} + S_{2y} \cdot S_{3y} + S_{2z} \cdot S_{3z}) \end{aligned}$$

And one can show using ladder operators (see Levine, Quantum Chem. (1991) p277)

$$S_x \alpha = +1/2 \beta \quad \text{and} \quad S_y \alpha = +1/2 \beta \quad \text{and} \quad S_x \beta = +1/2 \alpha \quad \text{and} \quad S_y \beta = -1/2 \alpha$$

using these results and applying the operators for the i^{th} spin ONLY to the i^{th} spin function gives

$$\Psi = \{\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\beta(3)\} / \sqrt{3}$$

$$Q1 \quad \hat{S}^2 = (\hat{S}_1 + \hat{S}_2 + \hat{S}_3) \cdot (\hat{S}_1 + \hat{S}_2 + \hat{S}_3) \quad \text{where } S_i \text{ act}$$

$$= \hat{S}_1^2 + \hat{S}_2^2 + \hat{S}_3^2 + 2 \cdot (S_1 \cdot S_2 + S_1 \cdot S_3 + S_2 \cdot S_3)$$

$$\text{and } S_i^1 S_j^2 = S_{ix}^1 |i\rangle S_{jx}^2 |j\rangle + S_{iy}^1 |i\rangle S_{jy}^2 |j\rangle + S_{iz}^1 |i\rangle S_{jz}^2 |j\rangle$$

$$\text{and } \begin{array}{lll} \hat{S}_x \alpha = \beta/2 & \hat{S}_x \beta = \alpha/2 & \hat{S}_z \alpha = \alpha/2 \\ \hat{S}_y \alpha = -i\beta/2 & \hat{S}_y \beta = -i\alpha/2 & \hat{S}_z \beta = -\beta/2 \end{array} \quad \left. \begin{array}{l} \text{here} \\ i = \sqrt{-1} \\ \hbar = 1 \end{array} \right.$$

$$\begin{array}{c} \hat{S} (\alpha \alpha \beta + \alpha \beta \alpha + \beta \alpha \alpha) \\ \{ 12 \text{ TERMS each} \\ 3 \text{ antisym} \} \end{array} \quad \begin{array}{c} \text{use position to identify spin index} \\ \text{① ② ③} \\ \alpha \alpha \beta \\ \alpha \beta \alpha \\ \beta \alpha \alpha \end{array} \quad \begin{array}{c} \text{TOTAL} \\ \text{① ② ③} \\ \beta \alpha \beta \\ \alpha \beta \beta \\ \beta \beta \beta \end{array}$$

\hat{S}_1^2	$\frac{3}{4} \alpha \alpha \beta$	$\frac{3}{4} \alpha \beta \alpha$	$\frac{3}{4} \beta \alpha \alpha$	$\frac{3}{4} \beta \beta \beta$
\hat{S}_2^2	$\frac{3}{4} \alpha \alpha \beta$	$\frac{3}{4} \alpha \beta \alpha$	$\frac{3}{4} \beta \alpha \alpha$	$\frac{3}{4} \beta \beta \beta$
\hat{S}_3^2	$\frac{3}{4} \alpha \alpha \beta$	$\frac{3}{4} \alpha \beta \alpha$	$\frac{3}{4} \beta \alpha \alpha$	$\frac{3}{4} \beta \beta \beta$

$S_{1x} S_{2x}$	$(\beta/2)(\beta/2)\beta$	$(\beta/2)(\alpha/2)\alpha$	$(\alpha/2)(\beta/2)\beta$	$\frac{1}{4}(\beta\beta\beta + \beta\alpha\alpha + \alpha\beta\beta)$
$S_{1y} S_{2y}$	$(i\beta/2)(-i\beta/2)\beta$	$(i\beta/2)(-i\alpha/2)\alpha$	$(-i\alpha/2)(i\beta/2)\beta$	$\frac{1}{4}(-\beta\beta\beta + \beta\alpha\alpha + \alpha\beta\beta)$
$S_{1z} S_{2z}$	$(\frac{1}{2})(\frac{1}{2})\beta$	$(\frac{1}{2})(\frac{1}{2})\alpha$	$(\beta/2)(\alpha/2)\beta$	$\frac{1}{4}\beta\beta\beta$

$S_{1x} S_{3x}$	$(\beta/2)\alpha(\alpha/2)$	$(\beta/2)\beta(\beta/2)$	$(\alpha/2)\alpha(\alpha/2)$	$\frac{1}{4}(\beta\alpha\alpha + \beta\beta\beta + \alpha\alpha\alpha)$
$S_{1y} S_{3y}$	$(i\beta/2)\alpha(-i\alpha/2)$	$(i\beta/2)\beta(-i\beta/2)$	$(-i\alpha/2)\alpha(-i\alpha/2)$	$\frac{1}{4}(\beta\alpha\alpha - \beta\beta\beta + \alpha\alpha\alpha)$
$S_{1z} S_{3z}$	$(\frac{1}{2})\alpha(\beta/2)$	$(\frac{1}{2})\beta(\beta/2)$	$(\beta/2)\alpha(\beta/2)$	$\frac{1}{4}\beta\beta\beta$

$S_{2x} S_{3x}$	$\alpha(\beta/2)(\alpha/2)$	$\alpha(\alpha/2)(\beta/2)$	$\beta(\beta/2)(\alpha/2)$	$\frac{1}{4}(\alpha\beta\alpha + \alpha\beta\beta + \beta\beta\beta)$
$S_{2y} S_{3y}$	$\alpha(-i\beta/2)(-i\alpha/2)$	$\alpha(-i\alpha/2)(i\beta/2)$	$\beta(i\beta/2)(-i\alpha/2)$	$\frac{1}{4}(-\alpha\beta\alpha - \alpha\beta\beta - \beta\beta\beta)$
$S_{2z} S_{3z}$	$\alpha(\frac{1}{2})(\beta/2)$	$\alpha(\beta/2)(\frac{1}{2})$	$\beta(\frac{1}{2})(\beta/2)$	$\frac{1}{4}\beta\beta\beta$

$$\hat{S}^2 \Psi = \frac{9}{4} \Psi + 2(\frac{3}{2}) \Psi = \frac{15}{4} \Psi \quad \Rightarrow \quad S = 3/2$$

Spin states of QUARTET $|\frac{3}{2}, \frac{1}{2}\rangle \quad (1s, m_s)$

The $(S_{ix} S_{ix} \text{ and } S_{iy} S_{iy})$ generate transform spin states that CANCEL

1a) What is the conserved component S_z ?

$S_z = s_{z1} + s_{z2} + s_{z3}$ where the small s_z act only on the spin of the i^{th} electron

so $S_z [(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3))/\sqrt{3}] =$

$$S_z = s_{z1}(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3))$$

$$+ s_{z2}(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3))$$

$$+ s_{z3}(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3))$$

$$= \left(\frac{1}{2}\right)\alpha(1)\alpha(2)\beta(3) + \left(\frac{1}{2}\right)\alpha(1)\beta(2)\alpha(3) + \left(\frac{-1}{2}\right)\beta(1)\alpha(2)\alpha(3)$$

$$+ \left(\frac{1}{2}\right)\alpha(1)\alpha(2)\beta(3) + \left(\frac{-1}{2}\right)\alpha(1)\beta(2)\alpha(3) + \left(\frac{1}{2}\right)\beta(1)\alpha(2)\alpha(3)$$

$$+ \left(\frac{-1}{2}\right)\alpha(1)\alpha(2)\beta(3) + \left(\frac{1}{2}\right)\alpha(1)\beta(2)\alpha(3) + \left(\frac{1}{2}\right)\beta(1)\alpha(2)\alpha(3)$$

$$= \left(\frac{1}{2}\right)(\alpha(1)\alpha(2)\beta(3) + \alpha(1)\beta(2)\alpha(3) + \beta(1)\alpha(2)\alpha(3)) = \left(\frac{1}{2}\right)^4 \Phi_2$$

1b) See additional page for solution of $S^2 \Phi_{1/2} = S(S+1) \Phi_{1/2} = (3/2)*(5/2) \Phi_{1/2} = 15/4 \Phi_{1/2}$

2. Of the atoms with $Z < 11$ which have ground states of odd parity?

② PARITY is symmetry w.r.t inversion

$$(x_i, y_i, z_i) \rightarrow (-x_i, -y_i, -z_i)$$

if $\pi_i \pi(i) = +1$ $\pi(i) = \text{state is even}$ (gerade)
 if $\pi_i \pi(i) = -1$ $\pi(i) = \text{state is odd}$ (ungerade)

For ATOMIC TERMS, the eigenvalue for $\sum \ell_i$

σ term with multiple unpaired spins is $(-1)^{\sum \ell_i}$

where ℓ_i is the orbital angular momentum of unpaired spin i

Thus if all even number of unpaired e^- 's \rightarrow EVEN

Symbol with \rightarrow ODD

H $1s$ $\sum \ell_i = 0$ g

He $1s^2$ g

Li $1s^2 2s$ $\sum \ell_i = 0$ g

Be $1s^2 2s^2$ g

B $2p$ $\sum \ell_i = 1$ u

C $2p^2$ $\sum \ell_i = 2$ g

N $2p^3$ $\sum \ell_i = 3$ u

O $2p^4$ $\sum \ell_i = 2$ g

F $2p^5$ $\sum \ell_i = 1$ u

Ne $2p^6$ $\sum \ell_i = 0$ g

3. Why is it incorrect to calculate the experimental ground state energy of lithium as $E_{2s} + 2^*E_{1s}$, where E_{2s} and E_{1s} are the experimental binding energies of the 1s and 2s electrons?

ANSWER: Because after removing the outermost 2s electron the remaining two electrons are more tightly bound. Similarly, after removing 2 electrons the last electron is more tightly bound.

E_{2s} is energy of $\text{Li } 1s^2 2s^1 \rightarrow \text{Li } 1s^2 2s^0 + e^-$ (~5 eV)
 E_{1s} is energy of $\text{Li } 1s^2 2s^1 \rightarrow \text{Li } 1s^1 2s^1 + e^-$ (~55 eV)

But true experimental ground state energy is energy for the process $\text{Li } 1s^2 2s^1 \rightarrow \text{Li}^{3+} + 3 e^-$

4a. Show that the commutation relations:

$[L_x, L_y] = i\hbar L_z$ with x,y,z cyclically permuted
 are equivalent to the single relationship $\hat{\mathbf{L}} \times \hat{\mathbf{L}} = i\hbar \hat{\mathbf{L}}$

$$\begin{aligned} \hat{\mathbf{L}} \times \hat{\mathbf{L}} &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ L_x & L_y & L_z \\ L_x & L_y & L_z \end{vmatrix} = (L_y L_z - L_z L_y) \hat{i} + (L_z L_x - L_x L_z) \hat{j} + (L_x L_y - L_y L_x) \hat{k} \\ &= [L_y, L_z] \hat{i} + [L_z, L_x] \hat{j} + [L_x, L_y] \hat{k} \end{aligned} \quad \begin{array}{l} \text{(has left off)} \\ \text{for clarity} \end{array}$$

Since $[L_x, L_y] = i\hbar L_z$ cyclically then

$$\hat{\mathbf{L}} \times \hat{\mathbf{L}} = i\hbar [L_x \hat{i} + L_y \hat{j} + L_z \hat{k}] = i\hbar \hat{\mathbf{L}}$$

4.b Evaluate $[L_x^2, L_y]$.

$$4) [L_x^2, L_y] = (L_x L_x L_y - L_y L_x L_x) \quad \xrightarrow{\text{evaluate from commutation}} \quad \text{evaluated from commutation}$$

$$\text{and } [L_x L_y] = i\hbar L_z = L_x L_y - L_y L_x$$

$$\text{so } L_x L_y = i\hbar L_z + L_y L_x \quad \text{and } L_y L_x = L_x L_y - i\hbar L_z$$

$$\begin{aligned} [L_x^2, L_y] &= \{ L_x (i\hbar L_z - L_y L_x) - (L_x L_y - i\hbar L_z) L_x \} \\ &= i\hbar L_x L_z + L_x L_y L_x - L_x L_y L_x + i\hbar L_z L_x \\ &= i\hbar (L_x L_z + L_z L_x) \end{aligned}$$

5. Show that $|n, t\rangle = e^{-iE_n t/\hbar} |n\rangle$ is a valid solution of the time dependent Schrödinger equation.

$$\hat{H}(t) |n, t\rangle = -\frac{\hbar}{i} \frac{\partial |n, t\rangle}{\partial t}$$

$$\hat{H}(t) e^{-iE_n t/\hbar} |n\rangle = -\frac{\hbar}{i} \frac{\partial}{\partial t} \left[e^{-iE_n t/\hbar} |n\rangle \right]$$

and $|n\rangle$ is not $f(t)$
 $e^{-iE_n t/\hbar}$ is a phase
so $\hat{H}(t)$ does not depend

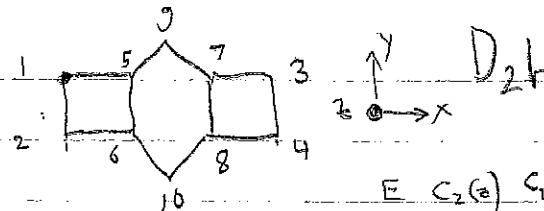
$$e^{-iE_n t/\hbar} \hat{H} |n\rangle = -\frac{\hbar}{i} |n\rangle \frac{\partial}{\partial t} (e^{-iE_n t/\hbar})$$

$$e^{-iE_n t/\hbar} \hat{H} |n\rangle = -\frac{\hbar}{i} |n\rangle \left(\frac{-iE_n}{\hbar} \right) e^{-iE_n t/\hbar}$$

$$\hat{H} |n\rangle = E_n |n\rangle \quad +, \text{independent S. E.}$$

$$\frac{\partial}{\partial x} e^{\alpha x} = \left(\frac{d\alpha}{dx} \right) e^{\alpha x}$$

⑥ Hückel
planar π



a) symmetry sets (1,2,3,4) A

(5,6,7,8) B

(9,10) C

$E \subset (e) \subset (4) \subset (8) \subset (8_{xy}) \subset (8_{xz}) \subset (8_{yz})$

$$\Gamma_A \quad 4 \quad 0 \quad 0 \quad 0 \quad 0 \quad -4 \quad 0 \quad 0$$

$$\Gamma_B \quad 4 \quad 0 \quad 0 \quad 0 \quad 0 \quad 0 \quad -4 \quad 0 \quad 0$$

$$\Gamma_C \quad 2 \quad 0 \quad -2 \quad 0 \quad 0 \quad -2 \quad 0 \quad 2$$

$$\begin{array}{ccccccc} & A_g & B_{1g} & B_{2g} & A_u & B_{3u} & B_{3u} \\ \Gamma_A = & 0 & 0 & 1 & 1 & 1 & 0 & 0 \end{array}$$

$$\Gamma_C = \begin{array}{ccccccc} & 0 & 0 & 0 & 1 & 0 & 1 & 0 & 0 \end{array}$$

b) SALS

$$B_{2g} (2): \chi_1 \phi_1 - \phi_4 - \phi_3 + \phi_2 - \phi_4 + \phi_1 + \phi_2 - \phi_3 = (\phi_1 - \phi_2 - \phi_3 - \phi_4)/2$$

$$\chi_2 (\phi_5 + \phi_6 - \phi_7 - \phi_8)/2$$

$$B_{3g} (3): \chi_3 \phi_1 - \phi_4 + \phi_3 - \phi_2 - \phi_4 + \phi_1 - \phi_2 + \phi_3 = \pm (\phi_1 - \phi_2 + \phi_3 - \phi_4)$$

$$\chi_4 \pm (\phi_5 - \phi_6 + \phi_7 - \phi_8)$$

$$\chi_5 \pm (\phi_9 - \phi_{10})$$

$$A_u (2): \chi_6 \phi_1 + \phi_4 = \phi_3 - \phi_2 + \phi_4 + \phi_1 - \phi_2 - \phi_3 = \pm \frac{1}{2} [\phi_1 - \phi_2 - \phi_3 + \phi_4]$$

$$\chi_7 \pm [\phi_5 - \phi_6 - \phi_7 + \phi_8]$$

$$B_{3u} (3): \chi_8 \phi_1 + \phi_4 + \phi_3 - \phi_2 + \phi_4 + \phi_1 + \phi_2 + \phi_3 = \pm \frac{1}{2} [\phi_1 + \phi_2 + \phi_3 + \phi_4]$$

$$\chi_9 \pm [\phi_5 - \phi_6 + \phi_7 + \phi_8]$$

$$\chi_{10} \pm (\phi_5 + \phi_{10})$$

c) Secular Determinant

B_{2g}

$$\begin{vmatrix} H_{11} - E & H_{12} \\ H_{21} & H_{22} - E \end{vmatrix} = \begin{vmatrix} \alpha + \beta - E & \beta \\ \beta & \alpha + \beta - E \end{vmatrix} = \begin{vmatrix} x+1 & 1 \\ 1 & x+1 \end{vmatrix} = (x+1)^2 - 1 = 0$$

$$H_{11} = \alpha$$

$$H_{12} = \beta$$

$$H_{11} = \frac{1}{4} (\phi_1 + \phi_2 - \phi_3 - \phi_4) H (\phi_1 + \phi_2 - \phi_3 - \phi_4) = \frac{1}{4} (\alpha + \beta + \beta + \alpha)$$

$$H_{12} = \frac{1}{4} (\phi_5 + \phi_6 - \phi_7 - \phi_8) = \alpha + \beta$$

$$\alpha = 6(\alpha - E)/4$$

$$\beta = \alpha - \beta x$$

$$H_{12} = \frac{1}{2} (\phi_1 + \phi_2 - \phi_3 - \phi_4) H (\phi_5 + \phi_6 - \phi_7 - \phi_8) = \frac{1}{4} (\beta + \beta + \beta + \beta) = \beta$$

$$x^2 + 2x = 0 \Rightarrow x(x+2) = 0 \quad \begin{cases} x=0 \\ x=-2 \end{cases}$$

$$\begin{array}{|c|} \hline E_1 = \alpha \\ \hline E_2 = \alpha + 2\beta \\ \hline \end{array}$$

A₁₁

$$\begin{vmatrix} H_{11} - E & H_{12} \\ H_{21} & H_{22} - E \end{vmatrix} = \begin{vmatrix} \alpha - \beta - E & \beta \\ \beta & \alpha - \beta - E \end{vmatrix} = \begin{vmatrix} x-1 & 1 \\ 1 & x-1 \end{vmatrix} = 0$$

$$H_{11} = \frac{1}{4}((\alpha_1 - \alpha_2 - \alpha_3 + \alpha_4)H(\alpha_1 - \alpha_2 - \alpha_3 + \alpha_4)) = \frac{1}{4}(\alpha - \beta + \beta + \alpha + \alpha - \beta - \beta + \alpha) = \alpha - \beta$$

$$H_{12} = \frac{1}{4}((\alpha_1 - \alpha_2 - \alpha_3 + \alpha_4)H(\alpha_5 - \alpha_6 - \alpha_7 + \alpha_8)) = \frac{1}{4}(\beta + \beta + \beta + \beta = \beta)$$

$$H_{21} = \frac{1}{4}((\alpha_5 - \alpha_6 - \alpha_7 + \alpha_8)H(\alpha_5 - \alpha_6 - \alpha_7 + \alpha_8)) = \frac{1}{4}(\alpha - \beta - \beta + \alpha + \alpha - \beta - \beta + \alpha) = \alpha - \beta$$

$$(x-1)^2 - 1 = x^2 - 2x + 1 - 1 = x^2 - 2x = 0 \Rightarrow x=0 \quad E_1 = \alpha$$

$$x(x-2) \quad x=+2 \quad E_2 = \alpha - 2\beta$$

B₃₃

$$\begin{vmatrix} H_{33} - E & H_{34} & H_{35} \\ H_{43} & H_{44} - E & H_{45} \\ H_{53} & H_{54} & H_{55} - E \end{vmatrix} = \begin{vmatrix} \alpha - \beta - E & \beta & 0 \\ \beta & \alpha - \beta - E & \sqrt{2}\beta \\ 0 & \sqrt{2}\beta & \alpha - E \end{vmatrix} = \begin{vmatrix} x-1 & 1 & 0 \\ 1 & x-1 & \sqrt{2} \\ 0 & \sqrt{2} & x \end{vmatrix} = 0$$

$$H_{33} = \frac{1}{4}(\alpha_1 - \alpha_2 + \alpha_3 - \alpha_4)H(\alpha_1 - \alpha_2 + \alpha_3 - \alpha_4) = \frac{1}{4}(\alpha - \beta + \alpha - \beta + \alpha - \beta + \alpha) = \alpha - \beta$$

$$H_{44} = \frac{1}{4}(\alpha_5 - \alpha_6 + \alpha_7 - \alpha_8)H(\alpha_5 - \alpha_6 + \alpha_7 - \alpha_8) = \frac{1}{4}(\alpha - \beta + \beta - \alpha + \alpha - \beta - \beta + \alpha) = \alpha - \beta$$

$$H_{55} = \frac{1}{2}(\alpha_9 - \alpha_{10})H(\alpha_9 - \alpha_{10}) = \frac{1}{2}(\alpha + \alpha) = \alpha$$

$$H_{34} = \frac{1}{4}(\alpha_1 - \alpha_2 + \alpha_3 - \alpha_4)H(\alpha_5 - \alpha_6 + \alpha_7 - \alpha_8) = \frac{1}{4}(\beta + \beta + \beta - \beta) = \beta$$

$$H_{35} = \frac{1}{2\sqrt{2}}(\alpha_1 - \alpha_2 + \alpha_3 - \alpha_4)H(\alpha_9 - \alpha_{10}) = 0$$

$$H_{45} = \frac{\sqrt{2}}{4}(\alpha_5 - \alpha_6 + \alpha_7 - \alpha_8)H(\alpha_9 - \alpha_{10}) = \beta + \beta + \beta - \beta = \sqrt{2}\beta \quad \text{NORM?}$$

$$(x-1) \begin{vmatrix} x-1 & \sqrt{2} \\ \sqrt{2} & x \end{vmatrix} - \begin{vmatrix} 1 & 0 \\ \sqrt{2} & x \end{vmatrix} = (x-1)(x(x-1) - 2) - x = 0$$

$$(x-1)(x^2 - x - 2) - x = x^3 - x^2 - 2x - x^2 + x + 2 - x = 0$$

$$x - x^3 - 2x^2 - 2x + 2 = 0$$

$$\text{Roots } \lambda_1 = -1.17 \quad E_3 = \alpha + 1.17\beta$$

$$\text{Wolfram } \lambda_2 = 0.69 \quad E_4 = \alpha - 0.69\beta$$

$$\lambda_3 = 2.48 \quad E_5 = \alpha - 2.48\beta$$

(B)

$$\begin{vmatrix} H_{88} - E & H_{89} & H_{8,10} \\ H_{98} & H_{99} - E & H_{9,10} \\ H_{10,8} & H_{10,9} & H_{10,10} - E \end{vmatrix} = \begin{vmatrix} \alpha + \beta - E & \beta & 0 \\ \beta & \alpha + \beta - E & \sqrt{2}\beta \\ 0 & \sqrt{2}\beta & \alpha - E \end{vmatrix} = \begin{vmatrix} x+1 & 1 & 0 \\ 1 & x+1 & \sqrt{2} \\ 0 & \sqrt{2} & x \end{vmatrix}$$

$$H_{88} = \frac{1}{4} (\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4) H (++++) = \frac{1}{4} (\alpha + \beta - \gamma) = \alpha + \beta$$

$$H_{99} = \frac{1}{4} (\alpha_5 + \alpha_6 + \alpha_7 + \alpha_8) H (L++++) = \frac{1}{4} (\alpha + \beta - \gamma) = \alpha + \beta$$

$$H_{10,10} = \frac{1}{2} (\alpha_9 + \alpha_{10}) H (\alpha_5 \alpha_{10}) = \frac{1}{2} (\alpha + \alpha) = \alpha$$

$$H_{85} = \frac{1}{4} (\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4) H (\alpha_5 + \alpha_6 + \alpha_7 + \alpha_8) = \frac{1}{2} (\beta + \beta + \beta + \beta) = \beta$$

$$H_{8,10} = \frac{1}{4} (\alpha_1 + \alpha_2 + \alpha_3 + \alpha_5) H (\alpha_9 + \alpha_{10}) = 0$$

$$H_{10,9} = \frac{\sqrt{2}}{24} (\alpha_5 + \alpha_6 + \alpha_7 + \alpha_8) H (\alpha_9 + \alpha_{10}) = \frac{\sqrt{2}}{2} (\beta + \beta + \beta + \beta) = \sqrt{2}\beta$$

$$(x+1) \begin{vmatrix} x+1 & \sqrt{2} \\ \sqrt{2} & x \end{vmatrix} = \begin{vmatrix} 1 & 0 \\ \sqrt{2} & x \end{vmatrix} = (x+1)[x(x+1) - 2] - x = 0$$

$$(x+1)(x^2 + x - 2) - x = x^3 + x^2 - 2x + x^2 + x - 2 - x = 0$$

$$x^3 + 2x^2 - 2x - 2 = 0$$

$$-2.48 \quad E_8 = \alpha + 2.48\beta$$

$$-0.69 \quad E_9 = \alpha + 0.69\beta$$

$$+1.17 \quad E_{10} = \alpha - 1.17\beta$$

$$1 - 1 - 1 \quad b_{2g} \text{ an}$$

$$1 - 1 - 1 + 1 \quad 1 - 1 - 1 + 1 \quad b_{3g}$$

$$1 - 1 - 1 + 1 \quad 1 - 1 + 1 - 1 \quad a_{1g} b_{3g}$$

$$1 + 1 - 1 - 1 + 1 \quad 1 - 1 \quad b_{2g} \text{ an}$$

$$\textcircled{e} \quad g_5 \quad b_{2g} a_u = {}^3 B_{2u} \quad \checkmark$$

$$\textcircled{f} \quad e_5$$

$$b_{3g} \rightarrow b_{3g} \rightarrow (a_u b_{3g}) \rightarrow {}^1 B_{3u}$$

$$\begin{array}{c} \uparrow \\ b_{2g} \end{array} \quad \begin{array}{c} \uparrow \\ a_u \end{array}$$

$$a_u \rightarrow b_{3g} \rightarrow (b_{2g}, b_{3g}) \rightarrow {}^1 B_{1g}$$

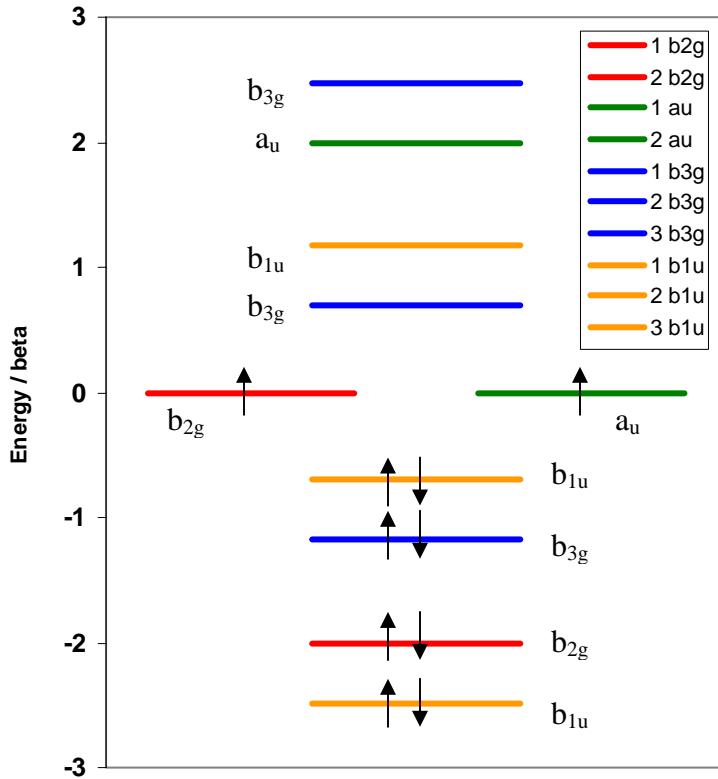
$${}^3 B_{2u} \rightarrow {}^3 B_{3u} \quad \Gamma_m = {}^1 B_{1g} \quad \text{forbidden}$$

$$\boxed{{}^3 B_{2u} \rightarrow {}^3 B_{3g}} \quad \Gamma_m = {}^1 B_{3u} \quad \text{allowed}$$

$\boxed{{}^3 B_{2u} \rightarrow {}^3 B_{3g}}$ can be VIBRONICALLY ALLOWED IF Γ_{1g}, b_{2u} or b_{3u} mode of vibrations is excited.

Ass#1 question #6

c) Quantitative energy level diagram for the pi system (Hückel, $\alpha = 0$, β is unit of E)



d) Lowest energy configuration is $(1b_{1u})^2 (1b_{2g})^2 (1b_{3g})^2 (2b_{1u})^2 (1b_{2g})^1 (1a_u)^1$

the corresponding STATES will have symmetries from $1b_{2g} \otimes 1a_u = B_{2u}$
 and the spins can be arranged either parallel (triplet) or antiparallel (singlet) $\rightarrow ^3B_{2u}$ and $^1B_{2u}$
 For states from the SAME configuration, the state with highest multiplicity is the lowest in energy, since this minimizes the electron repulsion (one of the Hund's rules)

Thus the **ground STATE** is $^3B_{2u}$

e) HOMO \rightarrow LUMO can be either $b_{2g} \rightarrow b_{3g}$ OR $a_u \rightarrow b_{3g}$
 $b_{2g} \rightarrow b_{3g} \rightarrow (a_u^1, b_{3g}^1)$ configuration has states $^3B_{3u}$ and $^1B_{3u}$
 $a_u \rightarrow b_{3g} \rightarrow (b_{2g}^1, b_{3g}^1)$ configuration has states $^3B_{1g}$ and $^1B_{1g}$

f) transitions

$^3B_{2u} \rightarrow ^3B_{3u}$ has a transition symmetry of $b_{2u} \otimes b_{3u} = B_{1g}$ FORBIDDEN

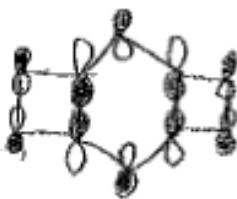
$^3B_{2u} \rightarrow ^3B_{1g}$ has a transition symmetry of $b_{2u} \otimes b_{1g} = B_{3u}$ ALLOWED

DIPOLE allowed transitions (i) conserve spin multiplicity; (ii) have $TM \subset b_{1u}, b_{2u}$ or b_{3u}
 Only transitions to the triplet states are allowed. $^3B_{2u} \rightarrow ^1B_{3u}$ and $^3B_{2u} \rightarrow ^1B_{1g}$ are FORBIDDEN

The $^3B_{2u} \rightarrow ^3B_{3u}$ can be made into an allowed vibronic transition if an odd number of quanta of vibrational modes of a_u, b_{2u} OR b_{3u} is created in the excited state.

(NB g-modes do not work; $b_{1g} \otimes b_{1u} = a_u$ to which (x,y,z) do NOT belong)

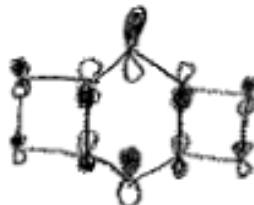
$3B_{2g}$ -2.4812



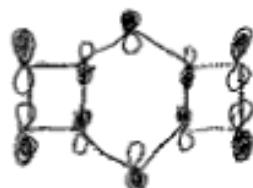
$2Au$ -2.0



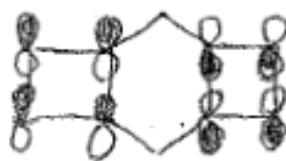
$3B_{2u}$ -1.1701



$2B_{2g}$ -0.6889

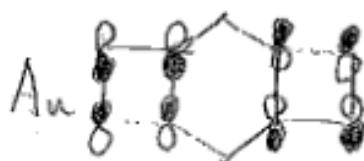


Ps_1

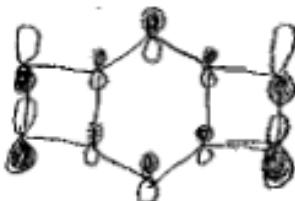


$2B_{3g}$ 0.0

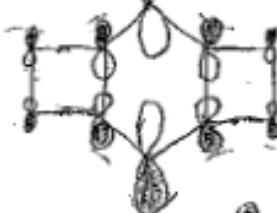
$1Au_g$ 0.0



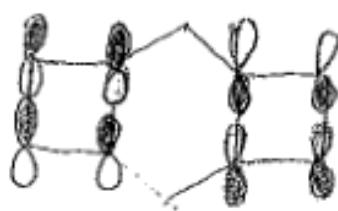
$2B_{2u}$ +0.6889



$1B_{2g}$ +1.1701



$1B_{3g}$ +2.0000



$1B_{2u}$ +2.4812

